Studies of track finding for long-lived particles at STCF*

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Reconstruction of the trajectories of charged particles at High Energy Physics experiments is a complicated task, in particular those of long-lived particles. At the future Super Tau-Charm Facility (STCF), long-lived particles are present in several important benchmark physics processes. A Common Tracking Software was used to reconstruct the trajectories of long-lived particles and it is shown that the track finding performance of the commonly used Combinatorial Kalman Filter for long-lived particles is limited by the seeding algorithm. This can be improved by steering the Combinatorial Kalman Filter with initial tracks provided by Hough Transform. The track finding performance of combined Hough Transform and Combinatorial Kalman Filter evaluated using the process $J/\psi \to \Lambda(\to p\pi^-)\bar{\Lambda}(\to \bar{p}\pi^+)$ at STCF is presented.

Keywords: Track finding, Common tracking software, Hough Transform, Long-lived particles

I. INTRODUCTION

Standard Model (SM) [1, 2] of particle physics includ-3 ing the unified Electro-Weak (EW) and Quantum Chromo-4 Dynamics (QCD) theories, has explained successfully almost 5 all experimental results about the microscopic world. How-6 ever, a couple of questions still remain, e.g. baryon asym-7 metry of the universe, dark matter, neutrino masses, num-8 ber of flavors. Beijing Electron Positron Collider (BEPCII) -9 Beijing Spectrometer (BESIII) [3] is the only multi-GeV $_{10}~e^{+}e^{-}$ collider operating in the au-charm sector, which provides an unique platform for studying non-perturbative QCD 12 and strong interactions of the SM. The Super Tau-Charm Fa-13 cility (STCF) [4, 5] is designed to continue and extend the 14 physics programs at BEPCII in near future, including probing 15 the nature of the strong interactions and hadron structure, pre-16 cise inspection of electroweak theories, exploring the asym-17 metry of matter-antimatter and searching for new physics be-18 yond the SM. STCF will operate at a center-of-mass-energy 19 of 2-7 GeV and a peak luminosity above 0.5×10^{35} cm⁻² s⁻¹, which is two orders higher than that at BEPCII.

The reconstruction of charged particles is the most fundamental and critical step in the data processing chain of High Energy Physics (HEP) experiments. To fulfill the physics goals and to further maximize the physics potential at the STCF, the charged particles need to be reconstructed with good efficiency. This includes not only those particles that decay immediately upon production but also the long-lived particles [6], e.g. the Λ and Ξ hyperons, which are relevant with a couple of important physics goals at STCF. For example, the weak decays of the Λ and Ξ hyperons provide promising channels for searching for new sources of CP violation [7–9].

 33 time-like nucleon and hyperon form factors for Q^2 values as 34 high as $40~{\rm GeV^2}$ [5]. Meanwhile, it is quite challenging to re- 35 construct the trajectories of long-lived particle decay products 36 because the long-lived particles may decay within or outside 37 the inner tracker hence having very limited number of hits 38 recorded by the inner tracker.

The Kalman Filter (KF) [10] algorithm is the most com-40 monly used algorithm for tracking in HEP and nuclear 41 physics. The Combinatorial Kalman Filter (CKF) [11, 12] is 42 an extended version of the KF, where the measurements are 43 progressively added to the track during the track propagation 44 steered by an initial estimate of the track parameters, i.e. seed. 45 The impact of magnetic field and material effects is incorpo-46 rated during track propagation hence CKF is capable to re-47 solve the hit ambiguity in a very dense tracking environment. 48 For this reason, CKF is deployed to find tracks by several 49 experiments e.g. ATLAS [13] and CMS [14], where thou-50 sands of tracks are present in a single event. CKF is also the 51 primary track finding algorithm at BelleII experiment [15]. 52 Recently, the CKF algorithm developed by BelleII experi-53 ment was reused to study the tracking performance [16] at 54 the Circular Electron-Positron Collider (CEPC) [17]. Despite 55 the great advantages of CKF, one downside of the KF-based 56 tracking algorithms is that they are subject to the performance 57 of the seeding algorithm, which might provide poor performance for long-lived particles. Recently, a track finding algo-59 rithm based on the Hough Transform used by BelleII experi-60 ment [18] and BESIII experiment [19] has been developed at 61 STCF [20], where tracking performance for prompt particles 62 without vertex displacement have been studied. It demonstrates promising tracking performance, in particular good ro-64 bustness against local hit inefficiency. However, the tracking 65 efficiency can be deteriorated by the presence of background 66 hits at low transverse momentum.

The ACTS (A Common Tracking Software) [21, 22] is an emerging open-source tracking software for HEP and nuclear physics experiments, with a suite of detector-agnostic and framework-independent modular track and vertex reconstruc-

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72 CKF algorithms in ACTS is underscored by their widespread 105 prises three layers of low-material budget silicon or gaseous ₇₃ adoption by experiments such as FASER [23], sPHENIX [24] ₁₀₆ detectors using either MAPS-based or μ -RWELL-based tech-₇₄ and a few R&D studies at STCF [25] and BESIII [26]. No- ₁₀₇ nology [28]. This study focuses on the μ -RWELL-based ITK 75 tably, ACTS has demonstrated its generality across a series 108 with the three layers placed at an inner radii of 60 mm, 110 76 of tracking detector types [27]. However, the performance of 109 mm, and 160 mm, respectively, and each layer has a thick-ACTS for long-lived particles has not been investigated. 77

78 ₇₉ fully gaseous tracking system consisting of a μ -RWELL [28]- ₁₁₂ in the z direction. For the MAPS-based ITK, the radii of the 80 based inner tracker and a drift chamber using combined 113 three silicon layers is 36 mm, 98 mm and 160 mm, respec-₈₁ Hough Transform and ACTS CKF to boost the tracking per- ₁₁₄ tively, and a hit resolution of 30 μ m \times 180 μ m is assumed. 82 formance for long-lived particles is studied. The manuscript 115 Unless explicitly specified, ITK denotes the μ -RWELL-based 83 is organized as follows. Section II presents a brief introduc- 116 ITK by default. 84 tion of the STCF detector. In Section III, the tracking work- 117 85 flow with different algorithms is introduced. Section IV fo- 118 tor, the Main Drift Chamber (MDC) operates using a 86 cuses on tracking performance for benchmark process with 119 He/C₃H₈(60/40) gas mixture and features a square cell con-87 long-lived particles at STCF. A brief conclusion is given in 120 figuration with a superlayer wire arrangement. The superlay-88 Section V.

II. STCF DETECTOR

The STCF detector [5] ensures comprehensive cover-91 age of the solid angle encompassing the collision point, depicted in Fig. 1. The STCF detector consists of a 93 tracking system comprising an Inner Tracker (ITK) and a 94 Main Drift Chamber (MDC), along with a Ring Imaging Cherenkov (RICH) detector [29] and a DIRC-like Time-of-96 Flight (DTOF) detector [30] for particle identification in the 129 97 barrel and endcap regions. Additionally, it incorporates a uniform Electro-magnetic Calorimeter (EMC) [31], a superconducting solenoid magnet generating 1 Tesla axial magnetic field, and a Muon Detector (MUD) positioned at the detector 101 system's outermost layer.

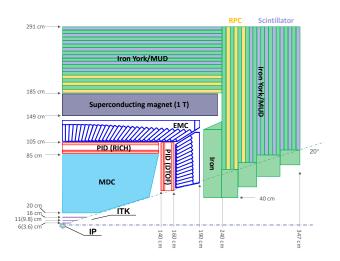


Fig. 1. Schematic layout of the STCF detector. The number in brackets indicate the radii of the MAPS-based ITK. Figure is taken from 151 Ref. [5].

To ensure optimal tracking efficiency for low-momentum 154 103 charged particles, the ITK within the tracking system cover- 155 transforming of experimental geometry, measurements, and

₇₁ tion algorithms. The promising performance of the KF and ₁₀₄ ing a polar angle range of 20° to 160° (i.e. $|\cos\theta| < 0.94$) com-110 ness of approximately 6.5 mm. It provides a spatial resolu-In this study, the tracking performance of STCF with a 111 tion around 100 μ m in the r- ϕ direction and around 400 μ m

> Central to the tracking system of the STCF detec-121 ers within the MDC alternate between stereo layers ("U" or "V") and axial layers, each containing six layers. In total, 123 the MDC comprises eight superlayers (AUVAUVAA) and 48 124 layers, with inner and outer radii of 200 mm and 850 mm, 125 respectively. The MDC provides spatial resolutions ranging between 120 μ m and 130 μ m.

TRACK RECONSTRUCTION USING COMBINED HOUGH TRANSFORM AND CKF

The workflow of track reconstruction using combined 130 Hough Transform and ACTS CKF is illustrated in Fig. 2. The 131 ACTS CKF is used to find track candidates through track fit-132 ting steered by the initial track parameters provided by either 133 ACTS seeding algorithm or Hough Transform algorithm de-134 veloped within the STCF offline software.

Interface between STCF offline software and ACTS

The Offline Software System of the Super Tau-Charm Fa-137 cility (OSCAR) [32, 33] serves as the offline event processing framework for the STCF. It provides common services for data processing and a suite of application tools dedicated to event generation, simulation, reconstruction, and physics analysis. For simulation purposes, OSCAR incorporates generation of τ -charm physics processes facilitated by the KKMC [34] generator, while particle decays are modeled with EvtGen as used by BESIII experiment [35], both seamlessly integrated within the framework. The STCF detector geometry is described using the Detector Description Toolkit, DD4Hep [36], with all geometric parameters stored in compact files utilizing the eXtensible Markup Language (XML) [37]. To simulate the interaction of particles with 150 the detector comprehensively, Geant4 [38] is integrated into OSCAR, ensuring a sophisticated full simulation. For track 152 reconstruction, the track finding algorithm based on Hough 153 Transform is developed in OSCAR.

The interface between OSCAR and ACTS facilitates the

157 tions. Geometry plugins within ACTS are tailored to stream- 191 verse momentum and the transverse impact parameter on the 158 line the conversion of experimental geometry representations, 192 x-y plane, which are required to satisfy the criteria optimized 159 such as DD4hep or TGeo [39], into ACTS internal geometry 193 for the relevant physics processes. The bending of the seed in 160 description. For ITK, the transformation involves convert- 194 the r-z plane is also required to be smaller than a threshold, cess entails transforming each sense wire within a drift cell 197 description of the ACTS seeding. 164 into a line surface. Leveraging dedicated material mapping 165 tools within ACTS, detailed material descriptions are pro-166 jected onto internal auxiliary surfaces of the ACTS geom-167 etry. For the conversion of measurements and initial track 168 parameters, two ROOT [40]-based readers have been developed. One reader extracts simulated hits from full simulation data and converts them into ACTS measurements taking into 171 account the resolution of the detectors. Another reader converts the initial estimate of the track parameters provided by 173 the Hough Transform algorithm into ACTS track parameters.

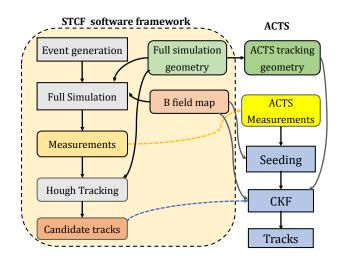


Fig. 2. The workflow of studying tracking performance using STCF software framework and ACTS.

ACTS seed finding

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The seeding algorithm in ACTS aims to find a few measurements which can provide position coordinates (x, y, z) in the global coordinate frame associated with a single particle to initiate the track following process. Without a seed, a particle cannot be reconstructed, hence the seed finding algorithm 180 acceptance region. 181

In a uniform magnetic field along the global z axis, the 183 helical trajectory of a charged particle is accurately defined by three measurements, thus forming a seed. In the case of STCF, these seeds are generated by combining three compat-186 ible measurements from the ITK detector with one measurement per ITK layer, as illustrated in Fig. 3. For each candidate seed, the curvature and center of the circle on the x-y plane 189 are determined using the Conformal Transform [41, 42]. Sub- 223, and for a circle with center (u, v) and radius r tangent to the

156 initial track estimates into corresponding ACTS representa- 190 sequently, these parameters are utilized to calculate the transing the signal readout unit tube within each μ -RWELL layer 195 which is optimized taking into account the impact of maginto sensitive cylinder surfaces. Similarly, for MDC, the pro- 196 netic field and multiple scattering. See Ref. [43] for a detailed

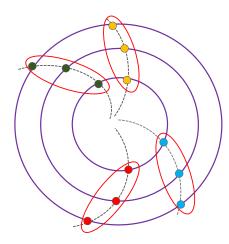


Fig. 3. Illustration of ACTS seeding using measurements from STCF ITK.

C. Track finding with Hough Transform

The principle of Hough Transform for track finding is illustrated in Fig. 4. With the presence of a magnetic field along global z axis, the projection of the track in the geometrical transverse x-y plane is a circle and the projection of the track in the geometrical s-z plane (s is the path length of the track in the x-y plane) is a straight line. The Conformal Transform can convert the projection of a track in the transverse x-y plane passing through the origin into a straight line, with a drift circle tangent to the projection of the track converted 208 to another circle tangent to the straight line, in the Conforu-v space. For a displaced track which has non-zero but 210 relatively small transverse impact track parameter, i.e. d_0 , 211 compared to the radius of the circle projected in the geomet- $_{212}$ rical transverse x-y plane, as shown in Fig. 4 left panel, the 213 trajectory of the transformed measurements (either a point or 214 a drift circle) on the Conformal space can be approximately 215 parameterized using a straight line, as shown in Fig. 4 midaims to find at least one seed for each particle in the detector 216 dle panel. The Hough Transform for track finding operates 217 on the principle that a straight line in the geometrical or Con-218 formal space can be described by two parameters, angle θ of $_{219}$ its normal and its algebraic distance ρ from the origin. For 220 a point with coordinate (u, v) on the straight line, it can be 221 transformed into a sinusoidal curve in the Hough space, i.e.

$$\rho = u \cdot \cos\alpha + v \cdot \sin\alpha \tag{1}$$

225 in the Hough space, i.e.

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$$\rho = u \cdot \cos\alpha + v \cdot \sin\alpha \pm r \tag{2}$$

The process of finding the measurements or drift circles $_{228}$ that arise from the same track in either the Conformal u-vspace or the geometrical s-z space becomes identifying the 230 curves that have an intersection in Hough space as shown in ²³¹ Fig. 4 right panel, i.e. a 2D histogram with optimized binning 232 taking into account the resolution of the track parameters are 283 tions. 233 filled if a curve passes through it and the peaks of the histograms are identified, as detailed in Ref. [20]. The parameters at the intersection can be converted to the track parameters describing the track projected to either the geometrical transverse x-y plane or the geometrical s-z plane. 237

The workflow of track finding using Hough Transform in 239 OCSCAR is shown in Fig.5. Initially, the measurements from ITK and MDC axial wires are used to reconstruct the projections of the tracks on the x-y plane, denoted as 2D tracks, 242 followed by circle fitting to extract track parameters of the 2D tracks. This is succeeded by associating the MDC stereo wire measurement candidates to the 2D tracks, where the z 245 position and path length s of the track at the stereo wires are 246 derived simultaneously. For each stereo wire measurement, two z position solutions can be obtained, and measurements from other tracks may be wrongly assigned to a 2D track. Therefore, a secondary application of Hough Transform is $_{250}$ employed to find the tracks in the s-z plane. More details can be found in Ref. [20].

Track finding with ACTS CKF

Starting from a set of initial track parameters, the ACTS 254 CKF is driven by the ACTS track propagator to search for 255 compatible measurements at a particular surface through KF 256 track fitting, as illustrated by Fig. 6. This process is also 257 known as track following. The measurement providing the 258 best fitting quality is associated to the track and used to filter 259 the track parameters for further track propagation.

PERFORMANCE STUDIES

Monte-Carlo samples

263 an important benchmark process at STCF allowing for several 318 layers particles traverse. Fig. 8 bottom panels shows the seedimportant physics studies relevant to Λ . Those signal events 319 ing efficiency as a function of particle p_T . It is observed that generated with the KKMC and EvtGen generators in OSCAR 320 Hough Transform algorithm can provide an efficiency above without considering bream-related backgrounds are used to 321 90% for $p(\bar{p})$ with p_T above 350 MeV/c and above 80% for π evaluate the tracking performance. The 2D distributions of 322 with p_T above 85 MeV/c, which is much improved compared the $\cos\theta$ versus p_T , vertex displacement in the x-y plane V_{xy} 323 to ACTS seeding. 269 versus p_T , and transverse impact track parameter d_0 versus 324 The reconstructed tracks are required to have at least five p_T , for proton (anti-proton), denoted as $p(\bar{p})$, and π in the 325 measurements on the track and have reconstructed $|\cos\theta|$ $_{271}$ $J/\psi \to \Lambda(\to p\pi^-)\bar{\Lambda}(\to \bar{p}\pi^+)$ events are shown in Fig. 7. 326 0.94. A reconstructed track is matched to its primary particle ₂₇₂ The π has a low momentum with p_T below 310 MeV/c and ₃₂₇ if the fraction of its hits from its primary particle, i.e. track

straight line, it can be transformed into two sinusoidal curves $_{273}$ $p(\bar{p})$ has a p_T up to 1.1 GeV/c. A non-negligible amount of 274 particles are decaying outside the first layer of ITK. However, $_{275}$ most of the tracks have d_0 smaller than the radius of the first (2) ₂₇₆ ITk layer, in particular for $p(\bar{p})$.

> Following event generation, Geant4 simulates hits from fi-278 nal state particles decaying from primary particles interacting 279 with the STCF tracking system in a uniform magnetic field 280 of 1T. Detector measurements are then generated by apply-281 ing Gaussian smearing to the positions of simulated hits, with 282 zero means and widths corresponding to the detector resolu-

Track finding performance

The performance of track finding, including seed finding 286 using either ACTS seeding algorithm or Hough Transform 287 algorithm at the first stage, and track following using ACTS 288 CKF at the second stage, is studied. Considering the accep-289 tance of STCF tracking system, only truth particles with p_T ₂₉₀ above 50 MeV and $|\cos\theta|$ below 0.94 are considered in the 291 performance metrics evaluation, which involves identifying 292 the primary particle [21] of a seed or a track, i.e. the simu-293 lated particle which has the most simulated hits contributing to this seed or track.

The seeding process serves as the initial step in track find-296 ing using CKF, which should provide seeds for all particles 297 in an ideal case. The ACTS seeding efficiency is defined as 298 the fraction of particles in the tracking system acceptance re-299 gion that have matched seeds with all three hits arising from 300 the same particle. The seeding efficiency using Hough Trans-301 form is defined by requiring that a matched seed has at least 50% hits from its primary particle.

For ACTS seeding, it's only possible to find seeds for a 304 track if the particle produces hits in all three layers of ITK, 305 indicating a vertex displacement below 66.5 mm. The com-306 parison between efficiencies of ACTS seeding and Hough Transform algorithm as a function of V_{xy} of the particles is 308 shown in Fig. 8 top panels. The ACTS seeding efficiency ap-309 proaches 100% when the number of measurements from ITK 310 is no less than 3. In particular, the ACTS seeding provides better seeding efficiency than Hough Transform for π with 312 small V_{xy} . However, ACTS seeding efficiency immediately 313 drops to zero if the number of measurements from ITK is below 3, indicating a significant limitation of ACTS seeding 315 algorithm, in particular for long-lived particles. The Hough Transform algorithm, functioning as a global tracking algo-The J/ψ decay process $J/\psi \to \Lambda(\to p\pi^-)\bar{\Lambda}(\to \bar{p}\pi^+)$ is 317 rithm, demonstrates reduced sensitivity to the number of ITK

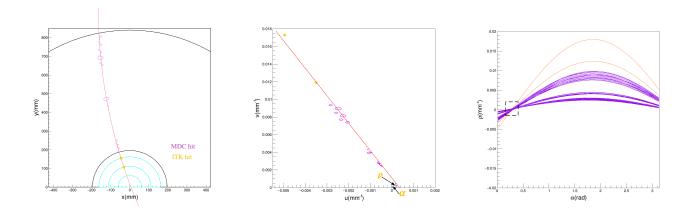


Fig. 4. An example of mapping detector measurements in geometrical transverse x-y plane (left) to the Conformal u-v space (middle) and eventually to the Hough curves in Hough space (right) for the particle p in a $J/\psi \to \Lambda(\to p\pi^-)\bar{\Lambda}(\to \bar{p}\pi^+)$ event. The particle decays between the first and the second layer of the ITK detector and hence has two hits on ITK. The hits or curves for ITK hits and MDC hits are shown by yellow and purple colors, respectively.

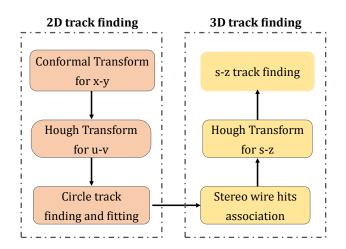


Fig. 5. The workflow of track finding using Hough Transform in OSCAR.

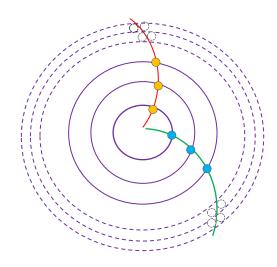


Fig. 6. Illustration of track finding using ACTS CKF with STCF ITK and MDC. Only two MDC layers are shown in the figure.

the track with the highest track purity is classified as the real 348 achieved using combined Hough Transform and CKF. track and others are classified as duplicate tracks. The track 349 among the reconstructed tracks. The duplicate rate is defined by the fraction of particles which have at least one duplicate 355 plicate rate than that using Hough Transform as seeding. track among the particles which have at least 5 simulated hits 356 in the detector acceptance region.

342 ticle V_{xy} and p_T . As expected, the tracking efficiency us- 359 ing efficiency, fake rate and duplicate rate of the combined $_{
m 343}$ ing ACTS seeding and CKF drops to zero when V_{xy} of the $_{
m 360}$ Hough Transform and ACTS CKF for the two different ITK 344 particle exceeds 66.5 mm, while the tracking efficiency with 361 designs. Since the first two layers of the MAPS-based

purity, is no less than 50%, and it's classified as a fake track 345 Hough Transform and CKF is less dependent on the particle if it's not matched to its primary particle. If more than one re- $_{346}$ V_{xy} . A tracking efficiency above 80% for $p(\bar{p})$ with p_T above constructed tracks are matched to the same simulated particle, $_{347}$ 350 MeV/c and above 70% for π with p_T above 85 MeV/c is

Figure 10 shows the fake rate and duplicate rate using the reconstruction efficiency is defined by the fraction of particles 350 two different seeding strategies. The fake rate is less than which have matched reconstructed tracks among the particles 351 0.4% and a non-negligible amount of duplicate tracks are which have at least 5 simulated hits in the detector acceptance 352 found for particles with p_T below 150 MeV/c, which have region. The fake rate is defined by the fraction of fake tracks 353 looping trajectories when traversing the detector in a mag-354 netic field. ACTS seeding results in lower fake rate and du-

The track finding performance of the alternative MAPS-357 based ITK for the long-lived particles is compared to that Figure 9 shows the tracking efficiency as a function of par- $_{358}$ of the μ -RWELL-based ITK. Figure 11 shows the track-

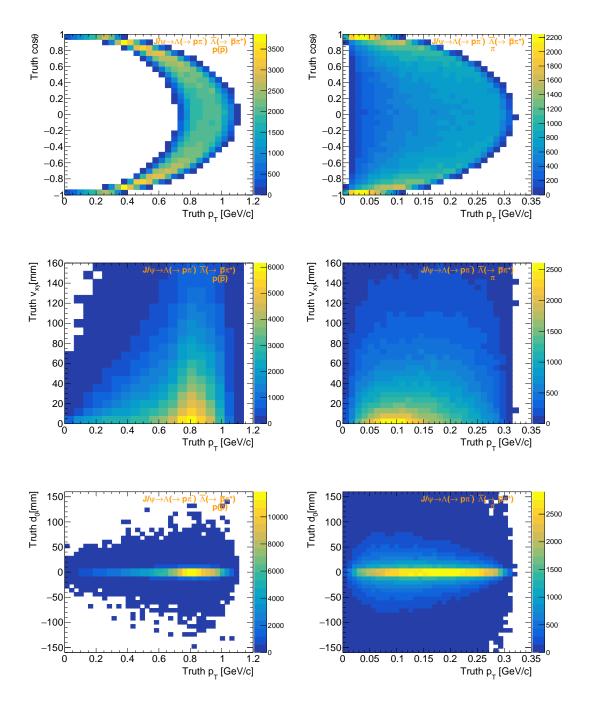


Fig. 7. The distributions of particle $\cos\theta$ versus p_T (top) and particle vertex displacement in the x-y plane V_{xy} versus p_T (middle), and d_0 versus p_T (bottom) for $p(\bar{p})$ (left) and π (right) in $J/\psi \to \Lambda(\to p\pi^-)\bar{\Lambda}(\to \bar{p}\pi^+)$ events.

 $_{362}$ ITK have smaller radii than the $\mu\text{-RWELL}\text{-based}$ ITK, the $_{367}$ MAPS-based ITK is less robust against the long-lived particles and provides slightly worse tracking efficiency than the $_{365}$ $\mu\text{-RWELL-based}$ ITK. The occurrence of fake and duplicate $_{368}$ tracks with the two designs is at similar level.

V. CONCLUSION

Processes with long-lived particles provide opportunities for probing CP, strong interaction etc. at the next generation Tau-Charm facility, STCF. However, high-performance track reconstruction for long-lived particles is a challenging and complicated task based on the tracking system of STCF. CKF is one of the most commonly used track find-

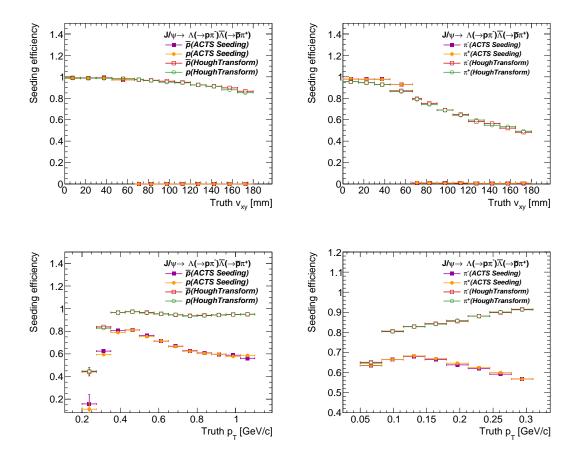


Fig. 8. The seeding efficiency as a function of the particle V_{xy} (top) and p_T (bottom) for $p(\bar{p})$ (left) and $\pi^+(\pi^-)$ (right) in 200k $J/\psi \to \Lambda(\to 0)$ $p\pi^-)\bar{\Lambda}(\to \bar{p}\pi^+)$ events. The solid purple square and yellow dot represent the results of ACTS seeding for particles with negative charge and positive charge, respectively. The hollow red square and green circle represent the results of Hough Transform for particles with negative charge and positive charge, respectively.

374 ing algorithms at HEP experiments with its performance sub- 390 85 MeV/c, with negligible occurrence of fake tracks. Dupli- $_{375}$ ject to the performance of the corresponding seeding algo- $_{391}$ cate tracks also exist, mainly arising from particles with p_T 376 rithm. For long-lived particles, CKF using traditional seeding 392 below 150 MeV/c with looping trajectories. Future devel-377 strategy, which often uses measurements from inner detec- 393 opment like extension of the 2D Hough space to 3D Hough $_{378}$ tor(s), demonstrates significant performance loss. Based on $_{394}$ space, where a track projection on the x-y plane not passing the STCF offline software and the common tracking software 395 through origin is described by three dedicated parameters, is ACTS, the combined performance of using Hough Trans- 396 foreseen to further enhance the tracking efficiency for longform as a seeding algorithm for ACTS CKF has been stud- 397 lived particles at STCF and beyond. 382 ied for the first time. The performance was evaluated us-383 ing $J/\psi \to \Lambda(\to p\pi^-)\bar{\Lambda}(\to \bar{p}\pi^+)$ events at STCF. The study shows that CKF steered by Hough Transform ends up with improved efficiency compared to CKF steered by traditional seeding algorithm for particles with large vertex displacement. The tracking efficiency using combined Hough 399 Transform and CKF is 80% for proton and anti-proton with 400 Science Foundation of China (NSFC) under Contract Nos. ₃₈₉ p_T above 350 MeV/c, and above 70% for π with p_T above ₄₀₁ 12375194, 12341504, 12375197, 12025502.

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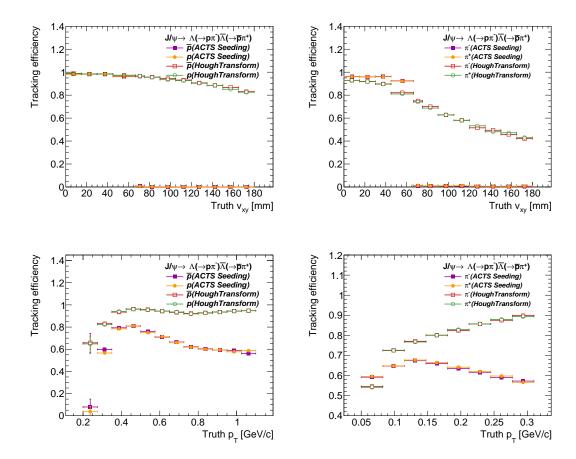


Fig. 9. The tracking efficiency as a function of the particle V_{xy} (top) and p_T (bottom) for $p(\bar{p})$ (left) and $\pi^+(\pi^-)$ (right) in 200k $J/\psi \to 0$ $\Lambda(\to p\pi^-)\bar{\Lambda}(\to \bar{p}\pi^+)$ events. The solid purple square and yellow dot represent the results with ACTS seeding for particles with negative charge and positive charge, respectively. The hollow red square and green circle represent the results with Hough Transform for particles with negative charge and positive charge, respectively.

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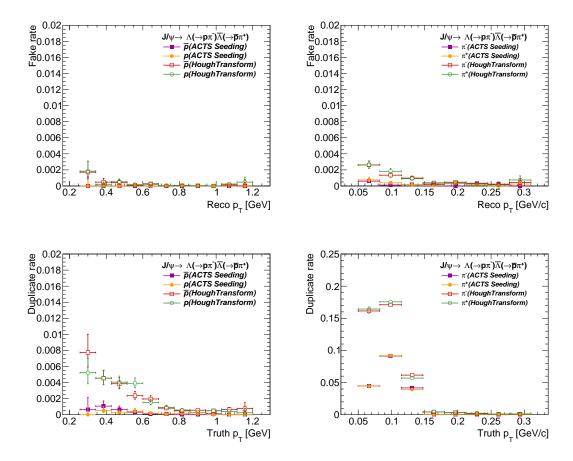


Fig. 10. The fake rate as a function of the track p_T (top) and duplicate rate as a function of the particle p_T (bottom) for $p(\bar{p})$ (left) and $\pi^+(\pi^-)$ (right) in 200k $J/\psi \to \Lambda(\to p\pi^-)\bar{\Lambda}(\to \bar{p}\pi^+)$ events. The solid purple square and yellow dot represent the results with ACTS seeding for particles with negative charge and positive charge, respectively. The hollow red square and green circle represent the results with Hough Transform for particles with negative charge and positive charge, respectively.

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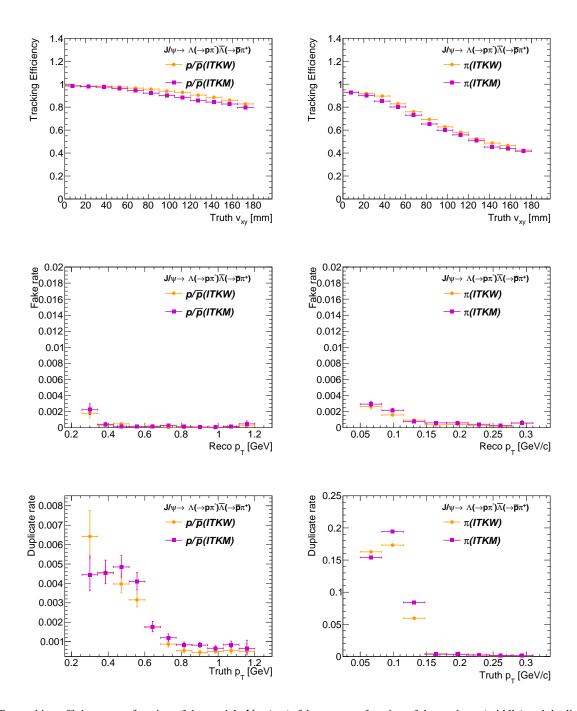


Fig. 11. The tracking efficiency as a function of the particle V_{xy} (top), fake rate as a function of the track p_T (middle) and duplicate rate as a function of the particle p_T (bottom) for $p(\bar{p})$ (left) and π (right) in 200k $J/\psi \to \Lambda(\to p\pi^-)\bar{\Lambda}(\to \bar{p}\pi^+)$ events using combined Hough Transform and ACTS CKF. The yellow dot and purple square represent the results with the μ -RWELL-based ITK (denoted as ITKW) and MAPS-based ITK (denoted as ITKM), respectively.

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